# Simulations in Statistical Physics

Course for MSc physics students

János Kertész

Lecture 11

Simulation of granular material (molecular dynamics = granular dynamics = **distinct element method**)

- Macroscopic particles (have temperature, pressure)
- Short range asymmetric interaction
- Dissipative collisions
- Friction

**Rigid particle model:** Angular momentum and friction has to be considered

Event driven



Perfect slip (no role of angular momentum)

Define "effective mass": 
$$m_{eff} \equiv \frac{m_i m_j}{m_i + m_j}$$

$$\overrightarrow{\Delta p_n} = -2m_{eff} \left[ \left( \vec{v}_i^{before} - \vec{v}_j^{before} \right) \vec{n} \right] \vec{n}$$

$$\vec{v}_i^{after} = \vec{v}_i^{before} - \frac{\overrightarrow{\Delta p_n}}{m_i} , \quad \vec{v}_j^{after} = \vec{v}_j^{before} + \frac{\overrightarrow{\Delta p_n}}{m_j}$$

For stick condition conservation of angular momentum has to be considered



Similarly, a coefficient for the tangential component can be introduced.

$$\begin{pmatrix} \mathbf{v}_{j}^{after} - \mathbf{v}_{i}^{after} \end{pmatrix} \cdot \mathbf{n} = -e_{n} \begin{pmatrix} \mathbf{v}_{j}^{before} - \mathbf{v}_{i}^{before} \end{pmatrix} \cdot \mathbf{n}$$
  

$$\Delta \mathbf{p} = -m_{eff} \left( e_{n} + 1 \right) \left[ \begin{pmatrix} \mathbf{v}_{j}^{before} - \mathbf{v}_{i}^{before} \end{pmatrix} \cdot \mathbf{n} \right] \cdot \mathbf{n}$$
 for perfect slip

Finite time singularity (inelastic collapse: numerical instability):

How long does it take until it stops?

$$t_{tot} = \sum_{j} t_{j} = 2\sqrt{\frac{2h^{initial}}{g}} \sum_{j} \sqrt{r^{j}} = 2\sqrt{\frac{2h^{initial}}{g}} \left(\frac{2}{1-\sqrt{r}} - 2\right)$$

 $\infty$  number of collisions in finite real time. Event driven simulation does not work.

http://www.haverford.edu/physics-astro/Gollub/vib\_granular/inelastic/inelastic.html

#### Ways out:

- Change *r* with the velocity (physics)  $r \sim (\Delta v)^{1/5}$
- Use soft potential and time steps
- Use different technique (Contact Dynamics)

### Soft potential with finite time steps

FViscoelastic model:



$$\varphi' = \theta(-x)(-kx - \gamma x)$$

Hertzian contact: Even material with linear properties will be nonlinear. For pure elasticity:

$$F_n^{elastic} = A\xi^{3/2}\theta(\xi)$$

Where A depends on the elastic moduli and the particle radius.

$$F_n^{diss} = B\sqrt{\xi}\,\dot{\xi}\theta(\xi)$$

B determines r

Similar, more complicated expressions are valid for  $F_{\perp}$ 



#### Shear

http://www.ph.biu.ac.il/~rapaport/java-apps/grshear.html

#### Vibrating plate

http://www.ph.biu.ac.il/~rapaport/java-apps/grvib.html



#### Algoritmikusan definiált modellek: Fraktál növekedés





Mineralization



### Surface crystallization



Elctrochem. deposition

Disordered viscous fingering

Bacterial colony growth



Basic equations:

- $\nabla^2 u = 0$  $\mathbf{v}\big|_{\Gamma} = -C\nabla u\big|_{\Gamma}$
- *u* scalar field (T, P, c...)
- $\mathbf{v}|_{\Gamma} = -C\nabla u|_{\Gamma}$  v velocity of the interface  $\Gamma$
- $u|_{\Gamma} = f(\nabla u, \kappa) \quad \kappa \text{ curvature (cutoff)}$

+ disorder



Laplacian or gradient governed groth:

If there is a bump, the gradient increases (c.f. electrostatic peak effect) the bump grows... instability

+ screening:

If 2 bumps grow, the faster will screen the slower and stop its growth

Simple model: Diffusion limited aggregation (DLA)

Start with a seed particle forming the initial aggregate.

\* Another particle comes from infinity via a random walk until it sticks to the aggregate.

Goto \*



100 million particles

Coarsened

Self-similar structure



1 million particles

In order to simulate (relatively) large samples tricks are needed

- No need to start from infinity: Birth ring sorrounding the aggregate

- No need to let the particles walk far away: killing ring
- If far from the aggregate: large steps possible

For very large (>10<sup>7</sup>) particles more tricks (fitted step size, dynamic storage)

Why are so terribly large aggregates needed?

Self similar fractals, scaling  $\rightarrow$  asymptotic behavior

How to measure fractal dimension?

## Dimensions

- Topological dimension:

Point:  $d_t=0$ , moving point:  $d_t=1$ , moving line,  $d_t=2...$ 

- Embedding dimension:

Number of independent directions

- Hausdorff (fractal) dimension



Area A is meaured by covering the object with squares of size  $\ell^2$ . # of such boxes:  $N_{\ell}$ .

$$A = \lim_{\ell \to 0} \ell^2 N_{\ell}$$



$$L = \lim_{\ell \to 0} \ell^1 N_{\ell}$$

In general:

$$M = \lim_{\ell \to 0} \ell^{d_t} N_{\ell} \qquad M: \text{ mass}$$

For a fractal this definition does not lead to a good result (0 or  $\infty$ )





$$\ell = \left(\frac{1}{3}\right)^n; \quad N_\ell = 4^n; \quad d_t = 1$$
$$M = \lim_{n \to \infty} \left(\frac{4}{3}\right)^n = \infty$$





Adapted from Benoit Mandelbrot, Fractals.

Instead of using  $d_t$  find the appropriate D fractal dimension such that

$$M = \lim_{\ell \to 0} \ell^{D} N_{\ell} \quad \text{is finite!}$$
  
Here:  $D = \frac{\ln 4}{\ln 3}$  (Def of fractal:  
 $d_{t} < D < d_{e}$ )

How to measure *D* for a random fractal?



There is always a lower and an upper cutoff (e.g., particle size, radius of gyration).

1. Box counting: Use the definition of *D*. Cover the object with a mesh of mesh size  $\ell$ , count the boxes where there is occupation. Plot log-log the dependence of  $N_{\ell} vs \ell$ .



2. Sand box method



 $M \sim L^D$ 

3. Correlation functions

$$C(r) = \langle \rho(r)\rho(0) \rangle \sim r^{-\alpha}$$
$$\int C(r')d^{d}r' \sim r^{r-\alpha} \sim M(r)$$
$$D = d - \alpha$$